

3. Fluctuation dissipation theorem

magnetization density $m(r)$

correlation function $\Gamma(r) = \langle m(r)m(0) \rangle - \langle m(r) \rangle \langle m(0) \rangle$

typical form $\Gamma(r) \sim r^{-p} e^{-r/\xi}$

fluctuation-dissipation theorem

For a translation invariant system we have

$$\chi = \frac{1}{V} \frac{\partial M}{\partial H} = \frac{1}{kT} \int d^3r \Gamma(r)$$

↑ magnetic susceptibility

Proof: $\mathcal{H} = \mathcal{H}_0 - \int H m(r) d^3r$

$$\chi = \frac{1}{V} \frac{\partial}{\partial H} \frac{1}{Z} \text{Tr} \int m(r) e^{-\beta \mathcal{H}} d^3r$$

$$= \frac{\beta}{V} \frac{1}{Z} \text{Tr} \int d^3r d^3r' m(r) m(r') e^{-\beta \mathcal{H}}$$

$$= \frac{\beta}{V} \frac{1}{Z} \text{Tr} \int d^3r m(r) e^{-\beta \mathcal{H}} \text{Tr} \int d^3r' m(r') e^{-\beta \mathcal{H}}$$

$$= \beta \int d^3r \langle m(r) m(0) \rangle - \beta \int d^3r \langle m(r) \rangle \langle m(0) \rangle$$

$$= \beta \int d^3r \Gamma(r)$$

4a) Critical Exponent

dimensionless temperature $t = \frac{T - T_c}{T_c}$

Field $h = \frac{H}{kT}$

heat capacity $C \sim t^{-\alpha}, T < T_c$
 $\sim t^{-\alpha'}, T > T_c$

magnetization $M \sim |t|^\beta, T < T_c$

susceptibility $\chi = \frac{1}{V} \partial_H M \sim |t|^{-\gamma}, T < T_c$

Equation of state $M \sim H^{1/\delta}, T = T_c$

Correlation function $\Gamma(r) \sim r^{-p} e^{-r/\xi}$ for $T \rightarrow T_c$

$$p = d - 2 + \nu$$
$$\xi \sim |t|^{-\nu}$$

	<u>Exp</u>	<u>Ising 3d</u>
α	0-0.14	0.12
β	0.32-0.39	0.31
γ	1.3-1.4	1.25
δ	4-5	5
ν	0.6-0.7	0.64
ν	0.05	0.05

There are relations between the critical exponents

e.g. Widom relation $\gamma = \beta(\delta - 1)$

exp $1.33 = 0.35 \times 3.5 = 1.23$

3d Ising $1.25 = 0.31 \times 4 = 1.24$

4b) Scaling

idea at the critical point there is only one scale: the correlation length ξ

$\Gamma(r) \sim r^{-(d-2+\eta)} = r^{-\nu/\xi}$

$\Gamma(r) = \langle m(r)m(0) \rangle - \langle m(r) \rangle \langle m(0) \rangle$

$\Rightarrow \int d^d r \Gamma(r) \sim M^2 \int d^d r r^{2-\eta} \Rightarrow \Gamma \sim \xi^{-\frac{d-2+\eta}{2}}$

$\chi \sim \int d^d r \Gamma(r) \Rightarrow \chi \sim \xi^{2-\eta}$

Free energy $\times \beta =$ dimensionless

$\Rightarrow \frac{F\beta}{V} \sim \xi^{-d}$

specific heat $C = T^2 \frac{\partial^2 F}{\partial T^2} \Rightarrow C \sim |\xi|^{-d+2}$

relations

$$\beta = \nu (d-2+n) k$$

$$\gamma = \nu (2-n)$$

Fisher

$$-2 = \nu d - 2$$

Josephson

(or hyperscaling because d is a parameter)

Next we derive a relation for δ

then $T \rightarrow T_c$ and H is small

in the partition function H occurs as

$$\begin{array}{l}
 \int d^d r \sim H M \\
 T \rightarrow T_c \quad M \sim \xi^{-(d-2+n)/2} \\
 d^d r \sim \xi^d
 \end{array}$$

to stay at constant physics we need $H \sim \xi^{\frac{-d}{2} + \frac{n}{2} - 1}$

but close to T_c $M \sim H^{\frac{1}{8}}$

$$\Rightarrow \xi^{-\frac{d-2+n}{2}} \sim \xi^{\frac{1}{8}(-\frac{d}{2} + \frac{n}{2} - 1)}$$

$$\Rightarrow \Rightarrow \beta \delta = \nu (2 + d - n) / 2$$

\Rightarrow 4 relations \Rightarrow we have only two independent exponents

The phase of the system also depends on two parameters, H and T . Is there a relation?